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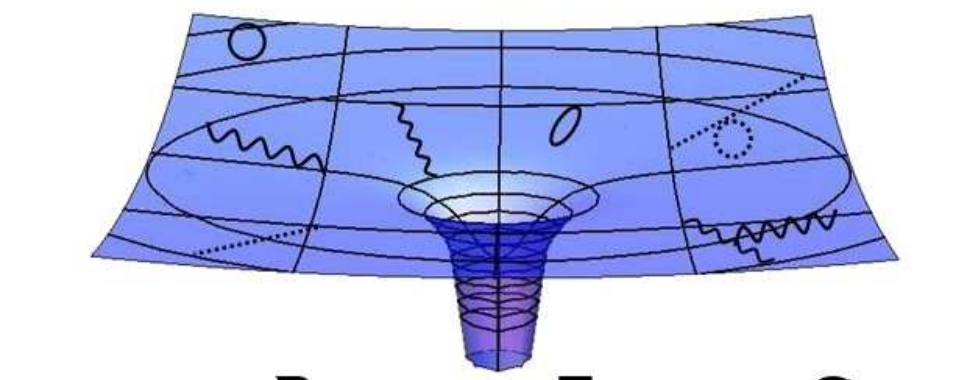
Spectral Study of Schrödinger Operators with Aperiodic Ordered Potentials

JOINT WORK WITH J. BELLISSARD AND G. DE NITTIS

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1. Motivation

We study Schrödinger operators induced by aperiodic solids (e.g. Quasicrystals). In the periodic case the Floquet-Bloch theory is a powerful tool to analyze the spectral properties. A similar tool for the aperiodic case is missing! Here, we investigate the approximation of non-periodic Schrödinger operators by periodic ones.

In general, it is impossible to approximate aperiodic Schrödinger operators

by periodic operators in the norm topology. To overcome this difficulty, in [2] we used weaker notions than the convergence in norm. This includes the convergence of the spectrum as well as the strong convergence of the operators. A detailed study of this convergence can be found in [3]. We use C^* -algebra techniques and provide an algorithmic construction of the approximants which will be helpful for concrete computations.

2. One-dimensional systems

Let \mathcal{A} be a finite set and $\mathcal{A}^{\mathbb{Z}} := \{\xi : \mathbb{Z} \rightarrow \mathcal{A}\}$ the set of bi-infinite words. Equipped with the product topology the space $\mathcal{A}^{\mathbb{Z}}$ is compact, Hausdorff and metrizable. The action $\tau : \mathcal{A}^{\mathbb{Z}} \rightarrow \mathcal{A}^{\mathbb{Z}}$ is defined by $(\tau\xi)(j) := \xi(j-1)$. A bi-infinite word $\xi \in \mathcal{A}^{\mathbb{Z}}$ is called *periodic* if there exists a $p \in \mathbb{N}$ such that $\xi(j+p) = \xi(j)$ for all $j \in \mathbb{Z}$.

We study closed, τ -invariant subsets $\Xi \subseteq \mathcal{A}^{\mathbb{Z}}$ which are called *subshifts*. The pair (Ξ, τ) is a topological dynamical system which represents a family of one-dimensional solids. For $\xi \in \mathcal{A}^{\mathbb{Z}}$ the *local pattern* $\mathcal{W}(\xi)$ is the set of all finite subwords of ξ . Then, $\mathcal{W}(\Xi)$ is the collection of all local patterns of elements in Ξ , i.e. $\mathcal{W}(\Xi) := \cup_{\xi \in \Xi} \mathcal{W}(\xi)$.

4. Main Result

Theorem. ([2]) Consider a family of self-adjoint Schrödinger operators $H_{\xi} : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$, $\xi \in \Xi$, associated to a periodically approximable subshift $\Xi \subseteq \mathcal{A}^{\mathbb{Z}}$.

Then, for $\xi \in \Xi$ there exists a sequence of periodic, self-adjoint Schrödinger operators $H_n : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$, $n \in \mathbb{N}$, such that:

- The sequence $(H_n)_{n \in \mathbb{N}}$ converges strongly to H_{ξ} for $n \rightarrow \infty$.
- The sequence of spectra converges, i.e. $\lim_{n \rightarrow \infty} \sigma(H_n) = \bigcup_{\xi \in \Xi} \sigma(H_{\xi})$ with respect to the Hausdorff metric on \mathbb{R} .
- If Ξ contains at least one non-periodic element the periods of $(H_n)_{n \in \mathbb{N}}$ grow to infinity for $n \rightarrow \infty$.
- The approximants are given explicitly by the closed paths that visit all vertices in the Gähler-Anderson-Putnam (De Bruijn) graphs.

3. Schrödinger Operator

For a subshift Ξ let $C^*(\Gamma)$ be the C^* -algebra associated to the groupoid $\Gamma := \Xi \rtimes_{\tau} \mathbb{Z}$. Each self-adjoint element $h \in C^*(\Gamma)$ defines a strongly continuous, covariant family of Schrödinger operators on $\ell^2(\mathbb{Z})$ by the left-regular representation $H_{\xi} := \pi_{\xi}(h) : \ell^2(\mathbb{Z}) \rightarrow \ell^2(\mathbb{Z})$, $\xi \in \Xi$. These operators are linear, bounded and self-adjoint.

A typical example for such a family of Schrödinger operators is

$$(H_{\xi}\psi)(j) := \psi(j-1) + \psi(j+1) + V(\tau^{-j}\xi)\psi(j), \quad \psi \in \ell^2(\mathbb{Z}), \quad (1)$$

where $V : \mathcal{A}^{\mathbb{Z}} \rightarrow \mathbb{R}$ is continuous and takes finitely many values (pattern equivariant function).

5. Periodically approximable

For a subshift Ξ a sequence of oriented, finite graphs $\mathcal{G}_n := (\mathcal{V}_n, \mathcal{E}_n, \partial_r, \partial_s)$, called the *Gähler-Anderson-Putnam* [1, 5] (De Bruijn [4]) *graphs*, is defined by

$$\begin{aligned} \mathcal{V}_n &:= \mathcal{W}(\Xi) \cap \mathcal{A}^n, & \text{(vertex set)} \\ \mathcal{E}_n &:= \mathcal{W}(\Xi) \cap \mathcal{A}^{n+1}, & \text{(edge set)} \\ \partial_s(a_0 \dots a_n) &:= a_0 \dots a_{n-1}, & \text{(source map)} \\ \partial_r(a_0 \dots a_n) &:= a_1 \dots a_n. & \text{(range map)} \end{aligned}$$

The n -th graph \mathcal{G}_n encodes the local structure of the system Ξ and the pos-

sible extensions of the patterns. The subshift is called *periodically approximable* if the graphs $(\mathcal{G}_n)_{n \in \mathbb{N}}$ are strongly connected, i.e. it exists a positively oriented closed path in $(\mathcal{G}_n)_{n \in \mathbb{N}}$ that visit all vertices.

Such paths define a sequence of periodic subshifts Ξ_n tending to Ξ as required in Box 6.

Lemma.(sufficient condition) If the dynamical system (Ξ, τ) is minimal, then, the subshift Ξ is periodically approximable.

6. Convergence of the spectrum

It is well-known that for a continuous self-adjoint section of a continuous field of C^* -algebras the spectrum behaves continuously. Since the Schrödinger operators associated to Ξ are elements of $C^*(\Xi \rtimes_{\tau} \mathbb{Z})$ we define a continuous field of C^* -algebras by a sequence of subshifts $(\Xi_n)_{n \in \mathbb{N}}$ tending to Ξ in the Hausdorff-topology. This convergence can be characterized by the convergence of the local patterns:

Lemma. A sequence of subshifts $(\Xi_n)_{n \in \mathbb{N}}$ converge to Ξ , if and only if the local patterns converge, i.e. for all $m \in \mathbb{N}$ exists $n_m \in \mathbb{N}$ such that $\mathcal{W}(\Xi_n) \cap \mathcal{A}^m = \mathcal{W}(\Xi) \cap \mathcal{A}^m$, $n \geq n_m$ where $\mathcal{A}^m := \{u : \{1, \dots, m\} \rightarrow \mathcal{A}\}$.

Let $(\Xi_n)_{n \in \mathbb{N}}$ be a sequence of subshifts tending to $\Xi_{\infty} := \Xi$ in the Hausdorff-topology. For a given family of Schrödinger operators $(H_{\xi})_{\xi \in \Xi}$ the approximating operators are chosen such that they form a continuous section on the related field of C^* -algebras $\pi : \prod_{n \in \mathbb{N}} C^*(\Xi_n \rtimes_{\tau} \mathbb{Z}) \rightarrow \overline{\mathbb{N}}$.

For instance, let $(H_{\xi})_{\xi \in \Xi}$ be defined as in Equation (1) in Box 3 and $\xi \in \Xi$. Then, for $\psi \in \ell^2(\mathbb{Z})$ the approximants are defined by

$$H_n\psi(j) := \psi(j-1) + \psi(j+1) + V(\tau^{-j}\eta)\psi(j),$$

where $\eta \in \Xi_n$ is chosen such that $\eta|_{N_n} = \xi|_{N_n}$ the restriction to the set $N_n := \{-\lfloor \frac{n}{2} \rfloor, \dots, \lfloor \frac{n}{2} \rfloor - 1\}$ around the origin. The last condition is needed to get the strong convergence of H_n to H_{ξ} .

7. The Fibonacci sequence

Let $\mathcal{A} := \{a, b\}$ and consider the Fibonacci substitution

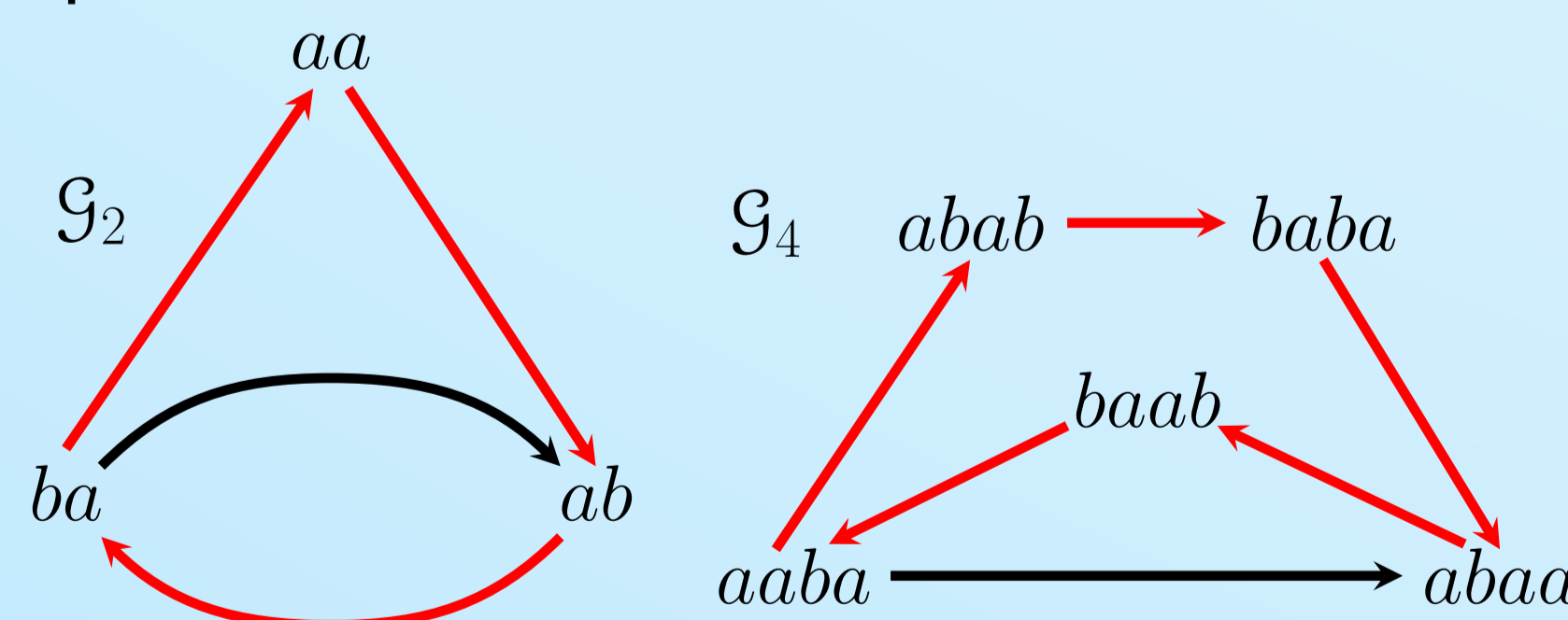
$$a \xrightarrow{S} ab, \quad b \xrightarrow{S} a.$$

Then, $\xi := \lim_{k \rightarrow \infty} S^{2k}(b) | S^{2k}(a)$ defines a fixed point of the substitution and

$$\xi = \dots ababaab|abaabab \dots \in \mathcal{A}^{\mathbb{Z}}.$$

The subshift $\Xi := \overline{\{\tau^k \xi \mid k \in \mathbb{Z}\}}$ is minimal and contains no periodic element. Hence, Ξ is periodically approximable and so the main result (Box 4.) applies to any family of Schrödinger operators $(H_{\xi})_{\xi \in \Xi}$ associated to the Fibonacci subshift.

The related Gähler-Anderson-Putnam graphs look as follows.



The red paths in \mathcal{G}_2 and \mathcal{G}_4 are closed paths that visit all vertices and creating the following periodic approximants

$$\Xi_2 : \dots aab|aab \dots \quad \Xi_4 : \dots abab|aabab|aaba \dots$$

The substitution rule can be used to compute the other Gähler-Anderson-Putnam graphs.

8. Future projects

(1) Fixing a metric on the space $\mathcal{A}^{\mathbb{Z}}$ the main result can be formulated as follows:

The Hausdorff convergence of subshifts $(\Xi_n)_{n \in \mathbb{N}}$ to Ξ implies the Hausdorff convergence of the spectrum.

Thus, it is natural to ask whether the rate of convergence of the spectrum can be estimated by the Hausdorff distance of Ξ_n to Ξ . In the case of finite range operators we expect an exponentially fast decay.

(2) So far the known techniques are only applicable to one-dimensional systems. Whereas our method could be extended to the higher dimensional case and so will provide a tool to study more complicated systems.

(3) In [3] the continuity of the spectrum is characterized by the continuity of the field of the corresponding C^* -algebra (as in our case). In this sense our result is optimal. In particular, if Ξ_n does not tend to Ξ then the spectrum of the associated Schrödinger operators will not converge in general. Thus, the creation of defects by breaking the convergence of the local patterns should be studied.

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